

VIVEKANANDA COLLEGE
THAKURPUKUR
KOLKATA-700063

NAAC ACCREDITED 'A' GRADE



Topic: System of Harmonic Oscillators (https://www.youtube.com/watch?v=1m7K_bmU61Y)

Course Title: Statistical Mechanics

Paper: PHY 423

Unit: 2

Semester: 2

Name of the Teacher: Arvind Pan

Name of the Department: Physics PG

Application of Canonical ensemble to two most important problems in physics:

1. Collection of Harmonic oscillators:
2. Statistics of Paramagnetism.

A system of Harmonic oscillators

N-independent, identical linear harmonic oscillators

$$\text{Hamiltonian } H = \sum_{i=1}^N \left(\frac{p_i^2}{2m} + \frac{1}{2} m \omega^2 q_i^2 \right)$$

→ Classical Approach

$$\int_{-\infty}^{+\infty} e^{-ax^2} dx = \sqrt{\frac{\pi}{a}}$$

$$\begin{aligned} Q_1 &= \frac{1}{h} \int e^{-\beta H} d^3 q_i d^3 p_i \\ &= \frac{1}{h} \int_{-\infty}^{+\infty} e^{-\beta p_i^2 / 2m} d^3 p_i \int_{-\infty}^{+\infty} e^{-\beta m \omega^2 q_i^2 / 2} d^3 q_i \\ &= \frac{1}{h} \sqrt{\frac{\pi}{\beta/2m}} \times \sqrt{\frac{\pi}{\beta m \omega^2 / 2}} = \frac{1}{\beta \hbar \omega} = \frac{kT}{\hbar \omega} \end{aligned}$$

Hence for the collection of N independent, distinguishable linear harmonic oscillators

$$Q = (Q_1)^N = (\beta \hbar \omega)^{-N}$$

No division by N! as these oscillators are representation of the energy levels, not particles or quasiparticles (photons & phonons) which are actually indistinguishable. Now the corresponding thermodynamic properties are

Helmholtz Free Energy $A = -kT \ln Q = N kT \ln (\hbar \omega / kT)$

Entropy $S = -\left(\frac{\partial A}{\partial T}\right)_{N,V} = -[Nk \ln (\hbar \omega / kT) + \frac{NkT}{\hbar \omega / kT} \left(-\frac{\hbar \omega}{kT^2}\right)]$

$$= Nk \ln \left(\frac{kT}{\hbar \omega}\right) + Nk$$

$$= Nk \left[\ln \left(\frac{kT}{\hbar \omega}\right) + 1 \right]$$

Internal Energy $U = A + TS = NkT \ln \left(\frac{\hbar \omega}{kT}\right) + NkT \left[\ln \left(\frac{kT}{\hbar \omega}\right) + 1 \right]$

$$= NkT$$

~~Specific heats~~ C_p

Pressure $P = -\left(\frac{\partial U}{\partial V}\right)_{N,T} = 0$

Chemical Potential $\mu = \left(\frac{\partial A}{\partial N}\right)_{V,T} = kT \ln \left(\frac{\hbar \omega}{kT}\right)$

~~Internal~~ Specific heats $C_p = C_v = \left(\frac{\partial U}{\partial T}\right)_{N,V} = Nk$

The last expression for internal energy / mean energy is in complete agreement with principle of equipartition of energy namely $2 \times \frac{1}{2} kT$ as we have two quadratic quadratic term in the Hamiltonian.

Let us do the above problem for quantum harmonic oscillator.

(b) Quantum mechanical approach: The energy for a quantum harmonic oscillator is:

$$E_n = \left(n + \frac{1}{2}\right) \hbar \omega, \quad n = 0, 1, 2, 3, \dots$$

Now single particle partition function is

$$Q_1 = \sum_{n=0}^{\infty} e^{-\beta E_n} = e^{-\beta \frac{\hbar \omega}{2}} \sum_{n=0}^{\infty} e^{-\beta n \hbar \omega}$$

$$= \frac{e^{-\beta \hbar \omega / 2}}{1 - e^{-\beta \hbar \omega}} = \frac{1}{e^{\beta \hbar \omega / 2} - e^{-\beta \hbar \omega / 2}}$$

$$= \frac{1}{2 \sinh(\beta \hbar \omega / 2)} = \left[2 \sinh\left(\frac{\beta \hbar \omega}{2}\right) \right]^{-1}$$

For N oscillator system

$$Q = (Q_1)^N = \left[2 \sinh\left(\frac{\beta \hbar \omega}{2}\right) \right]^{-N}$$

Thermodynamic Properties are

$$\begin{aligned} \# \text{ Helmholtz free Energy } A &= -kT \ln Q = NkT \ln \left[2 \sinh\left(\frac{\beta \hbar \omega}{2}\right) \right] \\ &= NkT \ln \left[e^{\beta \hbar \omega / 2} - e^{-\beta \hbar \omega / 2} \right] = NkT \ln \left[e^{\beta \hbar \omega / 2} (1 - e^{-\beta \hbar \omega}) \right] \\ &= N \left[kT \ln e^{\beta \hbar \omega / 2} + kT \ln (1 - e^{-\beta \hbar \omega}) \right] \end{aligned}$$

$$\begin{aligned} &= N \left[kT \beta \hbar \omega / 2 + kT \ln (1 - e^{-\beta \hbar \omega}) \right] \\ &= N \left[\frac{\hbar \omega}{2} + kT \ln (1 - e^{-\beta \hbar \omega}) \right] \end{aligned}$$

$$\# \text{ Entropy } S = -\left(\frac{\partial A}{\partial T}\right)_{N,V} = -Nk \ln \left[2 \sinh\left(\frac{\beta \hbar \omega}{2}\right) \right] - \frac{NkT}{2 \sinh\left(\frac{\beta \hbar \omega}{2}\right)} \coth\left(\frac{\beta \hbar \omega}{2}\right) \times \left(\frac{-\hbar \omega}{2kT^2}\right)$$

$$= Nk \frac{\beta \hbar \omega}{2} \coth\left(\frac{\beta \hbar \omega}{2}\right) - Nk \ln \left[2 \sinh\left(\frac{\beta \hbar \omega}{2}\right) \right] \left(\frac{-\hbar \omega}{2kT^2}\right)$$

$$= Nk \left[\frac{\beta \hbar \omega}{2} \frac{e^{\beta \hbar \omega / 2} + e^{-\beta \hbar \omega / 2}}{e^{\beta \hbar \omega / 2} - e^{-\beta \hbar \omega / 2}} - \ln \left[e^{\beta \hbar \omega / 2} (1 - e^{-\beta \hbar \omega}) \right] \right]$$

$$= Nk \left[\frac{\beta \hbar \omega}{2} \left(\frac{1 + e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right) - \frac{\beta \hbar \omega}{2} - \ln (1 - e^{-\beta \hbar \omega}) \right]$$

$$= Nk \left[\frac{\beta \hbar \omega}{2} \left(\frac{1 + e^{-\beta \hbar \omega} - 1 + e^{\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right) - \ln (1 - e^{-\beta \hbar \omega}) \right]$$

$$= Nk \left[\frac{\beta \hbar \omega}{2} \frac{2e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} - \ln (1 - e^{-\beta \hbar \omega}) \right]$$

$$= Nk \left[\frac{\beta \hbar \omega}{e^{\beta \hbar \omega} - 1} - \ln (1 - e^{-\beta \hbar \omega}) \right]$$

$$\# \text{ Internal Energy } U = A + TS = NkT \frac{\beta \hbar \omega}{2} \coth\left(\frac{\beta \hbar \omega}{2}\right)$$

$$= N \frac{\hbar \omega}{2} \left[\frac{e^{\beta \hbar \omega / 2} + e^{-\beta \hbar \omega / 2}}{e^{\beta \hbar \omega / 2} - e^{-\beta \hbar \omega / 2}} \right]$$

$$= \frac{N \hbar \omega}{2} \left[\frac{1 + e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} + 1 - 1 \right]$$

$$= \frac{N \hbar \omega}{2} \left[1 + \frac{1 + e^{-\beta \hbar \omega} - 1 + e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right]$$

$$= \frac{N \hbar \omega}{2} \left[1 + \frac{2e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} \right]$$

$$U = N \frac{\hbar \omega}{2} \left[1 + \frac{2}{e^{\beta \hbar \omega} - 1} \right] = N \left[\frac{\hbar \omega}{2} + \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1} \right]$$

Not in accordance with law of equipartition of energy, presence of zero point energy $\hbar \omega / 2$.

Pressure = $P = - \left(\frac{\partial U}{\partial V} \right)_{N, T} = 0$

Chemical potential $\mu = \frac{A}{N} = \left[\frac{\hbar \omega}{2} + kT \ln (1 - e^{-\beta \hbar \omega}) \right]$

Specific heat ($C_p = C_v$) = $\left(\frac{\partial U}{\partial T} \right)_{N, V} = N k \left(\frac{\hbar \omega}{2} \right) \left(\cos^2 \left(\frac{\beta \hbar \omega}{2} \right) \right) \times \left(\frac{-\hbar \omega}{2kT^2} \right)$

$$= N k \left(\frac{\hbar \omega}{2} \beta \right)^2 \cos^2 \left(\frac{\hbar \omega}{2} \beta \right)$$

$$= N k \left(\beta \hbar \omega \right)^2 \frac{e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$$

All these results would approach classical limits for $kT \gg \hbar \omega$

